# Gravity Control by means of *Electromagnetic Field* through *Gas* or *Plasma* at Ultra-Low Pressure

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It is shown that the gravity acceleration just above a chamber filled with gas or plasma at ultra-low pressure can be strongly reduced by applying an Extra Low-Frequency (ELF) electromagnetic field across the gas or the plasma. This Gravitational Shielding Effect is related to recent discovery of quantum correlation between gravitational mass and inertial mass. According to the theory samples hung above the gas or the plasma should exhibit a weight decrease when the frequency of the electromagnetic field is decreased or when the intensity of the electromagnetic field is increased. This Gravitational Shielding Effect is unprecedented in the literature and can not be understood in the framework of the General Relativity. From the technical point of view, there are several applications for this discovery; possibly it will change the paradigms of energy generation, transportation and telecommunications.

**Key words:** Phenomenology of quantum gravity, Experimental Tests of Gravitational Theories, Vacuum Chambers, Plasmas devices.

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#### I. INTRODUCTION

It will be shown that the local gravity acceleration can be controlled by means of a device called Gravity Control Cell (GCC) which is basically a recipient filled with gas or plasma where is applied an electromagnetic field. According to the theory samples hung above the gas or plasma should exhibit a weight decrease when the frequency of the electromagnetic field is decreased or when the intensity of the electromagnetic increased. The conductivity and the density of the gas or plasma are also highly relevant in this process.

With a GCC it is possible to convert the gravitational energy into rotational mechanical energy by means of the *Gravitational Motor*. In addition, a new concept of spacecraft (the *Gravitational Spacecraft*) and aerospace flight is presented here based on the possibility of gravity control. We will also see that the gravity control will be very important to *Telecommunication*.

#### II. THEORY

It was shown [1] that the relativistic gravitational mass  $M_g = m_g / \sqrt{1 - V^2/c^2}$ and the relativistic inertial mass  $M_i = m_{i0} / \sqrt{1 - V^2/c^2}$  are quantized, and given by  $M_{g} = n_{g}^{2} m_{i0(min)}$ ,  $M_{i} = n_{i}^{2} m_{i0(min)}$ where  $n_g$  and  $n_i$  are respectively, the gravitational quantum number and the inertial quantum number  $m_{i0(min)} = \pm 3.9 \times 10^{-73} kg$  is the elementary quantum of inertial mass. The masses  $m_{s}$  and  $m_{i0}$  are correlated by means of the following expression:

$$m_g = m_{i0} - 2 \sqrt{1 + \left(\frac{\Delta p}{m_i c}\right)^2} - 1 m_{i0}.$$
 (1)

Where  $\Delta p$  is the *momentum* variation on the particle and  $m_{i0}$  is the inertial mass at rest.

In general, the *momentum* variation  $\Delta p$  is expressed by  $\Delta p = F\Delta t$  where F is the applied force during a time interval  $\Delta t$ . Note that there is no restriction concerning the *nature* of the force F, i.e., it can be mechanical, electromagnetic, etc.

For example, we can look on the momentum variation  $\Delta p$  as due to absorption or emission of *electromagnetic energy* by the particle.

In the case of radiation,  $\Delta p$  can be obtained as follows: It is known that the radiation pressure, dP, upon an area dA = dxdy of a volume dV = dxdydz of a particle (the incident radiation normal to the surface dA) is equal to the energy dU absorbed per unit volume (dU|dV).i.e.,

$$dP = \frac{dU}{dV} = \frac{dU}{dxdydz} = \frac{dU}{dAdz}$$
 (2)

Substitution of dz = vdt (v is the speed of radiation) into the equation above gives

$$dP = \frac{dU}{dV} = \frac{\left(dU/dAdt\right)}{v} = \frac{dD}{v} \tag{3}$$

Since dPdA = dF we can write:

$$dFdt = \frac{dU}{v} \tag{4}$$

However we know that dF = dp/dt, then

$$dp = \frac{dU}{v} \tag{5}$$

From this equation it follows that

$$\Delta p = \frac{U}{v} \left(\frac{c}{c}\right) = \frac{U}{c} n_r$$

Substitution into Eq. (1) yields

$$m_g = \left\{ 1 - 2 \left[ \sqrt{1 + \left( \frac{U}{m_{i0}c^2} n_r \right)^2} - 1 \right] \right\} m_{i0}$$
 (6)

Where U, is the electromagnetic energy absorbed by the particle;  $n_r$  is the index of refraction.

Equation (6) can be rewritten in the following form

$$m_{g} = \left\{ 1 - 2 \left[ \sqrt{1 + \left( \frac{W}{\rho c^{2}} n_{r} \right)^{2}} - 1 \right] \right\} m_{i0}$$
 (7)

Where W = U/V is the density of electromagnetic energy and  $\rho = m_{i0}/V$  is the density of inertial mass.

The Eq. (7) is the expression of the quantum correlation between the gravitational mass and the inertial mass as a function of the density of electromagnetic energy. This is also the expression of correlation between gravitation and electromagnetism.

The density of electromagnetic energy in an *electromagnetic* field can be deduced from Maxwell's equations [2] and has the following expression

$$W = \frac{1}{2}\varepsilon E^2 + \frac{1}{2}\mu H^2 \tag{8}$$

It is known that  $B = \mu H$ ,  $E/B = \omega/k_r$  [3] and

$$v = \frac{dz}{dt} = \frac{\omega}{\kappa_r} = \frac{c}{\sqrt{\frac{\varepsilon_r \mu_r}{2} \left(\sqrt{1 + \left(\sigma/\omega\varepsilon\right)^2 + 1}\right)}}$$
 (9)

Where  $k_r$  is the real part of the propagation vector k (also called phase constant [4]);  $k = |\vec{k}| = k_r + ik_i$ ;  $\varepsilon$ ,  $\mu$  and  $\sigma$ , are the electromagnetic characteristics of the medium in which the incident (or emitted) radiation is propagating  $(\varepsilon = \varepsilon_r \varepsilon_0)$  where  $\varepsilon_r$  is the *relative* dielectric permittivity and  $\varepsilon_0 = 8.854 \times 10^{-12} F/m$ ;  $\mu = \mu_r \mu_0$  where  $\mu_r$  is the *relative* magnetic permeability and  $\mu_0 = 4\pi \times 10^7 H/m$ ;  $\sigma$  is the electrical conductivity). It is known that for *free-space*  $\sigma = 0$  and  $\varepsilon_r = \mu_r = 1$  then Eq. (9) gives

$$v = c \tag{10}$$

From (9) we see that the *index* of refraction  $n_r = c/v$  will be given by

$$n_r = \frac{c}{v} = \sqrt{\frac{\varepsilon_r \mu_r}{2} \left( \sqrt{1 + \left( \frac{\sigma}{\omega \varepsilon} \right)^2} + 1 \right)}$$
 (11)

Equation (9) shows that  $\omega/\kappa_r = v$ . Thus,  $E/B = \omega/k_r = v$ , i.e.,  $E = vB = v\mu H$ . Then, Eq. (8) can be rewritten in the following form:

$$W = \frac{1}{2} \left( \varepsilon \, v^2 \mu \right) \mu H^2 + \frac{1}{2} \, \mu H^2 \tag{12}$$

For  $\sigma \ll \omega \varepsilon$ , Eq. (9) reduces to

$$v = \frac{c}{\sqrt{\varepsilon_r \mu_r}}$$

Then, Eq. (12) gives

$$W = \frac{1}{2} \left[ \varepsilon \left( \frac{c^2}{\varepsilon_r \mu_r} \right) \mu \right] \mu H^2 + \frac{1}{2} \mu H^2 = \mu H^2$$
 (13)

This equation can be rewritten in the following forms:

$$W = \frac{B^2}{\mu} \tag{14}$$

or

$$W = \varepsilon E^2 \tag{15}$$

For  $\sigma \gg \omega \varepsilon$ , Eq. (9) gives

$$v = \sqrt{\frac{2\omega}{\mu\sigma}} \tag{16}$$

Then, from Eq. (12) we get

$$W = \frac{1}{2} \left[ \varepsilon \left( \frac{2\omega}{\mu \sigma} \right) \mu \right] \mu H^2 + \frac{1}{2} \mu H^2 = \left( \frac{\omega \varepsilon}{\sigma} \right) \mu H^2 + \frac{1}{2} \mu H^2 \cong$$

$$\cong \frac{1}{2} \mu H^2 \qquad (17)$$

Since  $E = vB = v\mu H$ , we can rewrite (17) in the following forms:

$$W \cong \frac{B^2}{2\mu} \tag{18}$$

or

$$W \cong \left(\frac{\sigma}{4\omega}\right) E^2 \tag{19}$$

By comparing equations (14) (15) (18) and (19) we see that Eq. (19) shows that the better way to obtain a strong value of W in practice is by applying an Extra Low-Frequency (ELF) electric field  $(w = 2\pi f \ll 1Hz)$  through a mean with high electrical conductivity.

Substitution of Eq. (19) into Eq. 7), gives

$$m_{g} = \left\{ 1 - 2 \sqrt{1 + \frac{\mu}{4c^{2}} \left(\frac{\sigma}{4\pi f}\right)^{3} \frac{E^{4}}{\rho^{2}}} - 1 \right\} m_{i0} \qquad (20)$$

This equation shows clearly that if an

electrical conductor mean has  $\rho <<1~Kg.m^{-3}$  and  $\sigma >>1$ , then it is possible obtain strong changes in its gravitational mass, with a relatively small ELF electric field. An electrical conductor mean with  $\rho <<1~Kg.m^{-3}$  is obviously a plasma.

There is a very simple way to test Eq. (20). It is known that inside a fluorescent lamp lit there is low-pressure Mercury plasma. Consider a 20W (80044 fluorescent lamp F20T12/C50/ECO GE, Ecolux® T12), whose characteristics and dimensions well-known [5]. Αt  $T \cong 318.15^{\circ} K$ , an optimum mercury vapor pressure of  $P = 6 \times 10^{-3} Torr = 0.8 N.m^{-2}$ is obtained, which is required for maintenance of high luminous efficacy throughout life. Under these conditions, the mass density of the Hg plasma can be calculated by means of the wellknown Equation of State

$$\rho = \frac{PM_0}{ZRT} \tag{21}$$

Where  $M_0 = 0.2006 kg.mol^{-1}$  is the molecular mass of the Hg;  $Z \cong 1$  is the compressibility factor for the Hg plasma;  $R = 8.314 joule.mol^{-1}.^0 K^{-1}$  is the gases universal constant. Thus we get

$$\rho_{Hg\ plasma} \cong 6.067 \times 10^{-5} \, kg.m^{-3}$$
(22)

The electrical conductivity of the Hg plasma can be deduced from the continuum form of Ohm's Law  $\vec{j} = \sigma \vec{E}$ , since the operating current through the lamp and the current density are well-known and respectively given by i=0.35A [5] and  $j_{lamp}=i/S=i/\frac{\pi}{4}\phi_{int}^2$ , where  $\phi_{int}=36.1mm$  is the inner diameter of the lamp. The voltage drop across the electrodes of the lamp is 57V [5] and the distance between them l=570mm. Then the electrical field along the lamp  $E_{lamp}$  is given by  $E_{lamp}=57V/0.570m=100~V.m^{-1}$ . Thus, we have

$$\sigma_{Hg\ plasma} = \frac{j_{lamp}}{E_{lamp}} = 3.419\ S.m^{-1}$$
 (23)

Substitution of (22) and (23) into (20) yields

$$\frac{m_{g(Hg\ plasm)}}{m_{i(Hg\ plasm)}} = \left\{1 - 2\sqrt{1 + 1.909 \times 10^{-17} \frac{E^4}{f^3}} - 1\right\} (24)$$

Thus, if an Extra Low-Frequency electric field  $E_{\rm ELF}$  with the following characteristics:  $E_{\rm ELF} \approx 100 V.m^{-1}$  and f < 1mHZ is applied through the Mercury plasma then a strong decrease in the gravitational mass of the Hg plasma will be produced.

It was shown [1] that there is an additional effect of *gravitational shielding* produced by a substance under these conditions. Above the substance the gravity acceleration  $g_1$  is reduced at the same ratio  $\chi = m_g / m_{i0}$ , i.e.,  $g_1 = \chi g$ , (g is the gravity acceleration *under* the substance). Therefore, due to the gravitational shielding effect produced by the decrease of  $m_{g(Hg\ plasma)}$  in the region where the ELF electric field  $E_{ELF}$  is applied, the gravity acceleration just above this region will be given by

$$g_{1} = \chi_{(Hg \ plasma)} g = \frac{m_{g(Hg \ plasma)}}{m_{i(Hg \ plasma)}} g =$$

$$= \left\{ 1 - 2 \left[ \sqrt{1 + 1.909 \times 10^{-17} \frac{E_{ELF}^{4}}{f_{ELF}^{3}}} - 1 \right] \right\} g \quad (25)$$

The trajectories of the electrons/ions through the lamp are determined by the electric field  $E_{lamp}$  along the lamp. If the ELF electric field across the lamp  $E_{ELF}$  is much greater than  $E_{lamp}$ , the current through the lamp can be interrupted. However, if  $E_{ELF} << E_{lamp}$ , these trajectories will be only slightly modified. Since here  $E_{lamp} = 100 \ V.m^{-1}$ , then we can arbitrarily choose  $E_{ELF}^{max} \cong 33 \ V.m^{-1}$ . This means that the maximum voltage drop, which can be applied across the metallic

(27)

plates, placed at distance d, is equal to the outer diameter (max  $^*$ ) of the bulb  $\phi_{lamp}^{max}$  of the 20W T-12 Fluorescent lamp, is given by

$$V_{max} = E_{ELF}^{max} \phi_{lamp}^{max} \cong 1.5 \ V$$

Since  $\phi_{lamp}^{max} = 40.3mm$  [5].

Substitution of  $E_{ELF}^{max} \cong 33 \ V.m^{-1}$  into (25) yields

$$g_{1} = \chi_{(Hg \ plasma)}g = \frac{m_{g(Hg \ plasma)}}{m_{i(Hg \ plasma)}}g = \left\{1 - 2\left[\sqrt{1 + \frac{2.264 \times 10^{-11}}{f_{ELF}^{3}}} - 1\right]\right\}g$$
 (26)

Note that, for  $f < 1mHz = 10^{-3}Hz$ , the gravity acceleration can be strongly reduced. These conclusions show that the ELF Voltage Source of the set-up shown in Fig.1 should have the following characteristics:

- Voltage range: 0 - 1.5 V

- Frequency range:  $10^{-4}$ Hz  $- 10^{-3}$ Hz

In the experimental arrangement shown in Fig.1, an ELF electric field with  $E_{ELF} = V/d$ intensity crosses fluorescent lamp; V is the voltage drop plates of the across the metallic  $d = \phi_{lamp}^{max} = 40.3mm.$ capacitor and When the ELF electric field is applied, the gravity acceleration just above the lamp (inside the dotted box) decreases according to (25) and the changes can be measured by means of the system balance/sphere presented on the top of Figure 1.

In Fig. 2 is presented an experimental arrangement with *two* fluorescent lamps in order to test the gravity acceleration above the *second* lamp. Since gravity acceleration above the *first* lamp is given by  $\vec{g}_1 = \chi_{1(Hg\ plasma)}\vec{g}$ , where

 $\chi_{1(Hg plasma)} = \frac{m_{g1(Hg plasma)}}{m_{i1(Hg plasma)}} =$   $= \left\{ 1 - 2 \sqrt{1 + 1.909 \times 10^{-17} \frac{E_{ELF(1)}^4}{f_{ELF(1)}^3} - 1} \right\}$ 

Then, above the *second* lamp, the gravity acceleration becomes

$$\vec{g}_2 = \chi_{2(Hg\ plasma)}\vec{g}_1 = \chi_{2(Hg\ plasma)}\chi_{1(Hg\ plasma)}\vec{g}$$
 (28) where

$$\chi_{2(Hg\ plasma)} = \frac{m_{g\,2(Hg\ plasma)}}{m_{i2(Hg\ plasma)}} = \left\{ 1 - 2 \left[ \sqrt{1 + 1.909 \times 10^{-17} \frac{E_{ELF(2)}^4}{f_{ELF(2)}^3} - 1} \right] \right\}$$
(29)

Then, results

$$\frac{g_2}{g} = \left\{ 1 - 2 \left[ \sqrt{1 + 1.909 \times 10^{-17} \frac{E_{ELF(1)}^4}{f_{ELF(1)}^3}} - 1 \right] \right\} \times \left\{ 1 - 2 \left[ \sqrt{1 + 1.909 \times 10^{-17} \frac{E_{ELF(2)}^4}{f_{ELF(2)}^3}} - 1 \right] \right\} \tag{30}$$

From Eq. (28), we then conclude that if  $\chi_{1(Hg\ plasma)} < 0$  and also  $\chi_{2(Hg\ plasma)} < 0$ , then  $g_2$  will have the same direction of g. This way it is possible to intensify several times the gravity in the direction of  $\vec{g}$ . On the other hand, if  $\chi_{1(Hg\ plasma)} < 0$  and  $\chi_{2(Hg\ plasma)} > 0$  the direction of  $\vec{g}$  will be contrary to direction of  $\vec{g}$ . In this case will be possible to intensify and become  $\vec{g}_2$  repulsive in respect to  $\vec{g}$ .

If we put a lamp above the *second* lamp, the gravity acceleration above the *third* lamp becomes

$$\vec{g}_3 = \chi_{3(Hg \ plasma)} \vec{g}_2 =$$

$$= \chi_{3(Hg \ plasma)} \chi_{2(Hg \ plasma)} \chi_{1(Hg \ plasma)} \vec{g} \quad (31)$$

or

<sup>\*</sup> After heating.

$$\frac{g_3}{g} = \left\{ 1 - 2 \left[ \sqrt{1 + 1.909 \times 10^{-17} \frac{E_{ELP(1)}^4}{f_{ELP(1)}^3}} - 1 \right] \right\} \times \left\{ 1 - 2 \left[ \sqrt{1 + 1.909 \times 10^{-17} \frac{E_{ELP(2)}^4}{f_{ELP(2)}^3}} - 1 \right] \right\} \times \left\{ 1 - 2 \left[ \sqrt{1 + 1.909 \times 10^{-17} \frac{E_{ELP(3)}^4}{f_{ELP(3)}^3}} - 1 \right] \right\} \times \left\{ 1 - 2 \left[ \sqrt{1 + 1.909 \times 10^{-17} \frac{E_{ELP(3)}^4}{f_{ELP(3)}^3}} - 1 \right] \right\} \tag{32}$$

If  $f_{ELF(1)} = f_{ELF(2)} = f_{ELF(3)} = f$  and

$$E_{\it ELF(1)} = E_{\it ELF(2)} = E_{\it ELF(3)} = V/\phi =$$
 
$$= V_0 \sin \omega t/40.3mm =$$
 
$$= 24.814 V_0 \sin 2\pi ft.$$

Then, for t = T/4 we get

$$E_{ELF(1)} = E_{ELF(2)} = E_{ELF(3)} = 24.814V_0$$
.

Thus, Eq. (32) gives

$$\frac{g_3}{g} = \left\{ 1 - 2 \left[ \sqrt{1 + 7.237 \times 10^{-12} \frac{V_0^4}{f^3}} - 1 \right] \right\}^3$$
 (33)

For  $V_0 = 1.5V$  and f = 0.2mHz (t = T/4 = 1250s = 20.83min) the gravity acceleration  $\vec{g}_3$  above the *third* lamp will be given by

$$\vec{g}_3 = -5.126\vec{g}$$

Above the *second* lamp, the gravity acceleration given by (30), is

$$\vec{g}_2 = +2.972\vec{g}$$

According to (27) the gravity acceleration above the *first* lamp is

$$\vec{g}_1 = -1,724\vec{g}$$

Note that, by this process an acceleration  $\vec{g}$  can be increased several times in the direction of  $\vec{g}$  or in the opposite direction.

In the experiment proposed in Fig. 1, we can start with ELF voltage sinusoidal wave of amplitude  $V_0=1.0V$  and frequency 1mHz. Next, the frequency will be progressively decreased down to 0.8mHz, 0.6mHz, 0.4mHz and 0.2mHz. Afterwards, the amplitude of the voltage wave must be increased to

 $V_0 = 1.5V$  and the frequency decreased in the above mentioned sequence.

Table1 presents the *theoretical* values for  $g_1$  and  $g_2$ , calculated respectively by means of (25) and (30). They are also plotted on Figures 5, 6 and 7 as a function of the frequency  $f_{\it ELF}$ .

Now consider a chamber filled with Air at  $3\times 10^{-12}torr$  and 300K as shown in Figure 8 (a). Under these circumstances, the mass density of the air inside the chamber, according to Eq. (21) is  $\rho_{air} \cong 4.94\times 10^{-15}\,kg.m^{-3}$ .

If the frequency of the *magnetic* field, B, through the air is f=60Hz then  $\omega\varepsilon=2\pi f\varepsilon\cong 3\times 10^{-9}\,S\,/\,m$ . Assuming that the electric conductivity of the air inside the chamber,  $\sigma_{(air)}$  is much less than  $\omega\varepsilon$ , i.e.,  $\sigma_{(air)}<<\omega\varepsilon$  (The atmospheric air conductivity is of the order of  $2-100\times 10^{-15}\,S.m^{-1}$  [6, 7]) then we can rewritten the Eq. (11) as follows

$$n_{r(air)} \cong \sqrt{\varepsilon_r \mu_r} \cong 1$$
 (34)

From Eqs. (7), (14) and (34) we thus obtain

$$m_{g(air)} = \left\{ 1 - 2 \left[ \sqrt{1 + \left( \frac{B^2}{\mu_{air} \rho_{air} c^2} n_{r(air)} \right)^2} - 1 \right] \right\} m_{i(air)} =$$

$$= \left\{ 1 - 2 \left[ \sqrt{1 + 3.2 \times 10^6 B^4} - 1 \right] m_{i(air)} \right\}$$
(35)

Therefore, due to the *gravitational* shielding effect produced by the decreasing of  $m_{g(air)}$ , the gravity acceleration above the *air* inside the chamber will be given by

$$g' = \chi_{air} g = \frac{m_{g(air)}}{m_{i(air)}} g =$$

$$= \left\{ 1 - 2 \left[ \sqrt{1 + 3.2 \times 10^6 B^4} - 1 \right] \right\} g$$

Note that the gravity acceleration above the *air* becomes *negative* for  $B > 2.5 \times 10^{-2} T$ .

For B = 0.1T the gravity acceleration above the air becomes

$$g' \cong -32.8g$$

Therefore the ultra-low pressure air inside the chamber, such as the Hg plasma inside the fluorescent lamp, works like a Gravitational Shield that in practice, may be used to build *Gravity Control Cells* (GCC) for several practical applications.

Consider for example the GCCs of Plasma presented in Fig.3. The ionization of the plasma can be made of several manners. For example, by means of an electric field between the electrodes (Fig. 3(a)) or by means of a RF signal (Fig. 3(b)). In the first case the ELF electric field and the ionizing electric field can be the same.

Figure 3(c) shows a GCC filled with air (at ambient temperature and 1 atm) strongly ionized by means of alpha particles emitted from 36 radioactive ions sources (a very small quantity of Americium 241 $^{\dagger}$ ). The radioactive element Americium has a half-life of 432 years, and emits alpha particles and low energy gamma rays ( $\approx 60 KeV$ ). In order to shield the alpha particles and gamma rays emitted from the Americium 241 it is

sufficient to encapsulate the GCC with epoxy. The alpha particles generated by the americium ionize the oxygen and nitrogen atoms of the air in the ionization chamber (See Fig. 3(c)) increasing the electrical conductivity of the air inside the chamber. The high-speed alpha particles hit molecules in the air and knock off electrons to form ions, according to the following expressions

$$O_2 + H_e^{++} \rightarrow O_2^+ + e^- + H_e^{++}$$
  
 $N_2 + H_e^{++} \rightarrow N_2^+ + e^- + H_e^{++}$ 

It is known that the electrical conductivity is proportional to both the concentration and the mobility of the *ions* and the *free electrons*, and is expressed by

$$\sigma = \rho_e \mu_e + \rho_i \mu_i$$

Where  $\rho_e$  and  $\rho_i$  express respectively the concentrations  $\left(C/m^3\right)$  of electrons and ions;  $\mu_e$  and  $\mu_i$  are respectively the mobilities of the electrons and the ions.

In order to calculate the electrical conductivity of the air inside the ionization chamber, we first need to calculate the concentrations  $\rho_e$  and  $\rho_i$ . We start calculating the *disintegration constant*,  $\lambda$ , for the Am 241:

$$\lambda = \frac{0.693}{T^{\frac{1}{2}}} = \frac{0.693}{432(3.15 \times 10^7 \, \text{s})} = 5.1 \times 10^{-11} \, \text{s}^{-1}$$

Where  $T^{\frac{1}{2}} = 432 years$  is the half-life of the Am 241.

One kmole of an isotope has mass equal to atomic mass of the isotope expressed in kilograms. Therefore, 1g of Am 241 has

$$\frac{10^{-3} kg}{241 \ kg/kmole} = 4.15 \times 10^{-6} \, kmoles$$

One kmole of any isotope contains the Avogadro's number of atoms. Therefore 1g of Am 241 has

$$N = 4.15 \times 10^{-6} \text{ kmoles} \times$$
$$\times 6.025 \times 10^{26} \text{ atoms/kmole} = 2.50 \times 10^{21} \text{ atoms/kmole}$$

 $<sup>^{\</sup>dagger}$  The radioactive element *Americium* (Am-241) is widely used in ionization smoke detectors. This type of smoke detector is more common because it is inexpensive and better at detecting the smaller amounts of smoke produced by flaming fires. Inside an ionization detector there is a small amount (perhaps 1/5000th of a gram) of americium-241. The Americium is present in oxide form (AmO2) in the detector. The cost of the AmO<sub>2</sub> is US\$ 1,500 per gram. The amount of radiation in a smoke detector is extremely small. It is also predominantly alpha radiation. Alpha radiation cannot penetrate a sheet of paper, and it is blocked by several centimeters of air. The americium in the smoke detector could only pose a danger if inhaled.

Thus, the activity [8] of the sample is

$$R = \lambda N = 1.3 \times 10^{11}$$
 disintegrations/s.

However, we will use 36 ionization sources each one with 1/5000th of a gram of Am 241. Therefore we will only use  $7.2 \times 10^{-3} g$  of Am 241. Thus, R reduces to:

$$R = \lambda N \cong 10^9$$
 disintegrations/s

This means that at *one* second, about  $10^9 \, \alpha \ particles$  hit molecules in the air and knock off electrons to form ions  $O_2^+$  and  $N_2^+$  inside the ionization chamber. Assuming that *each* alpha particle yields *one* ion at each  $1/10^9$  second then the total number of ions produced in one second will be  $N_i \cong 10^{18} ions$ . This corresponds to an ions concentration

$$\rho_i = eN_i/V \approx 0.1 /V \quad (C/m^3)$$

Where V is the volume of the ionization chamber. Obviously, the concentration of electrons will be the same, i.e.,  $\rho_e=\rho_i$ . For d=2cm and  $\phi=20cm$  (See Fig.3(c)) we obtain

 $V = \frac{\pi}{4} (0.20)^2 (2 \times 10^{-2}) = 6.28 \times 10^{-4} m^3$  The n we get:

$$\rho_a = \rho_i \approx 10^2 \ C/m^3$$

This corresponds to the minimum concentration level in the case of conducting materials. For materials, at temperature of 300K, the mobilities  $\mu_e$  and  $\mu_i$  vary from 10 to  $100 \ m^2 V^{-1} s^{-1}$  [9]. Then we can assume  $\mu_e = \mu_i \approx 10 \ m^2 V^{-1} s^{-1}$ . (minimum mobility level for conducting materials). Under these conditions, the electrical conductivity of the air inside ionization chamber is

$$\sigma_{air} = \rho_e \mu_e + \rho_i \mu_i \approx 10^3 \, \text{S.m}^{-1}$$

At temperature of 300K, the air density inside the GCC, is  $\rho_{air}=1.1452kg.m^{-3}$ . Thus, for d=2cm,  $\sigma_{air}\approx 10^3\,S.m^{-1}$  and f=60Hz Eq. (20) gives

$$\chi_{air} = \frac{m_{g(air)}}{m_{i(air)}} =$$

$$= \left\{ 1 - 2 \left[ \sqrt{1 + \frac{\mu}{4c^2} \left( \frac{\sigma_{air}}{4\pi f} \right)^3 \frac{V_{rms}^4}{d^4 \rho_{air}^2}} - 1 \right] \right\} =$$

$$= \left\{ 1 - 2 \sqrt{1 + 3.10 \times 10^{-16} V_{rms}^4} - 1 \right\}$$

Note that, for  $V_{rms}\cong 7.96KV$ , we obtain:  $\chi_{(air)}\cong 0$ . Therefore, if the voltages range of this GCC is: 0-10KV then it is possible to reach  $\chi_{air}\cong -1$  when  $V_{rms}\cong 10KV$ .

It is interesting to note that  $\sigma_{air}$  can be strongly increased by increasing the amount of Am 241. For example, by using 0.1g of Am 241 the value of R increases to:

$$R = \lambda N \cong 10^{10}$$
 disintegrations/s

This means  $N_i \cong 10^{20} ions$  that yield

$$\rho_i = eN_i/V \approx 10 /V (C/m^3)$$

Then, by reducing, d and  $\phi$  respectively, to 5mm and to 11.5cm, the volume of the ionization chamber reduces to:

$$V = \frac{\pi}{4} (0.115)^2 (5 \times 10^{-3}) = 5.19 \times 10^{-5} m^3$$
  
Consequently, we get:

$$\rho_e = \rho_i \approx 10^5 \, C/m^3$$

Assuming that  $\mu_e = \mu_i \approx 10 \ m^2 V^{-1} s^{-1}$ , then the electrical conductivity of the air inside the ionization chamber becomes

$$\sigma_{air} = \rho_e \mu_e + \rho_i \mu_i \approx 10^6 \, S.m^{-1}$$

This reduces for  $V_{rms} \cong 18.8V$  the voltage necessary to yield  $\chi_{(air)} \cong 0$  and reduces

to  $V_{rms} \cong 23.5V$  the voltage necessary to reach  $\chi_{air} \cong -1$ .

If the outer surface of a metallic sphere with radius a is covered with a radioactive element (for example Am 241), then the electrical conductivity of the air (very close to the sphere) can be strongly increased (for example up to  $\sigma_{air} \cong 10^6 \, s.m^{-1}$ ). By applying a low-frequency electrical potential  $V_{ms}$  to the sphere, in order to produce an electric field  $E_{rms}$  starting from the outer surface of the sphere, then very close to the sphere the low-frequency electromagnetic field is  $E_{rms} = V_{rms}/a$ , and according to Eq. (20), the  $gravitational\ mass$  of the air in this region expressed by

$$m_{g(air)} = \left\{ 1 - 2 \sqrt{1 + \frac{\mu_0}{4c^2} \left( \frac{\sigma_{air}}{4\pi f} \right)^3 \frac{V_{rms}^4}{a^4 \rho_{air}^2} - 1} \right\} m_{i0(air)},$$

can be easily reduced, making possible to produce a controlled *Gravitational Shielding* (similar to a GCC) surround the sphere.

This becomes possible to build a spacecraft to work with a gravitational shielding as shown in Fig. 4.

The *gravity accelerations* on the spacecraft (due to the rest of the Universe. See Fig.4) is given by

$$g'_i = \chi_{air} g_i$$
  $i = 1, 2, 3 ... n$ 

Where  $\chi_{air}=m_{g(air)}/m_{i0(air)}$  . Thus, the gravitational forces acting on the spacecraft are given by

$$F_{is} = M_g g'_i = M_g (\chi_{air} g_i)$$

By reducing the value of  $\chi_{air}$ , these forces can be reduced.

According to the *Mach's principle*;

"The *local inertial forces* are determined by the *gravitational interactions* of the local system with the distribution of the cosmic masses".

Thus, the local inertia is just the gravitational influence of the rest of the existing in Universe. Consequently, if we reduce the gravitational interactions between spacecraft and the rest of the Universe, then the inertial properties of the spacecraft will be also reduced. This effect leads to a new concept of spacecraft and space flight.

Since  $\chi_{air}$  is given by

$$\chi_{air} = \frac{m_{g(air)}}{m_{i0(air)}} = \left\{ 1 - 2 \left[ \sqrt{1 + \frac{\mu_0}{4c^2} \left( \frac{\sigma_{air}}{4\pi f} \right)^3 \frac{V_{rms}^4}{a^4 \rho_{air}^2}} - 1 \right] \right\}$$

Then, for  $\sigma_{air} \cong 10^6 \, s.m^{-1}$ , f = 6Hz, a = 5m,  $\rho_{air} \cong 1Kg.m^{-3}$  and  $V_{rms} = 3.35KV$  we get

$$\chi_{air} \cong 0$$

Under these conditions, the gravitational forces upon the spacecraft become approximately nulls and consequently, the spacecraft practically *loses its inertial properties*.

Out of the terrestrial atmosphere, the gravity acceleration upon the spacecraft is negligible and therefore the gravitational shielding is not necessary. However, if the spacecraft is in the outer space and we want to use the gravitational shielding then,  $\chi_{air}$  must be replaced by  $\chi_{vac}$  where

$$\chi_{vac} = \frac{m_{g(vac)}}{m_{i0(vac)}} = \left\{ 1 - 2 \left[ \sqrt{1 + \frac{\mu_0}{4c^2} \left( \frac{\sigma_{vac}}{4\pi f} \right)^3 \frac{V_{rms}^4}{a^4 \rho_{vac}^2} - 1} \right] \right\}$$

The electrical conductivity of the ionized outer space (very close to the spacecraft) is small; however, its density is remarkably small (<<  $10^{-16}\,Kg.m^{-3}$ ), in such a manner that the smaller value of the factor  $\sigma_{vac}^3/\rho_{vac}^2$  can be easily compensated by the increase of  $V_{rms}$ .

It was shown that, when the gravitational mass of a particle gravitational mass reduced to the between  $+0.159M_{\odot}$ ranging  $-0.159M_i$ , it becomes imaginary [1], i.e., the gravitational and the inertial masses particle become imaginary. Consequently, the particle disappears from our ordinary space-time. However, factor  $\chi = M_{g(imaginary)}/M_{i(imaginary)}$ remains real because

$$\chi = \frac{M_{g(imaginary)}}{M_{i(imaginary)}} = \frac{M_{g}i}{M_{i}i} = \frac{M_{g}}{M_{i}} = real$$

Thus, if the gravitational mass of the particle is reduced by means of absorption of an amount of electromagnetic energy U, for example, we have

$$\chi = \frac{M_g}{M_i} = \left\{ 1 - 2 \left[ \sqrt{1 + \left( U / m_{i0} c^2 \right)^2} - 1 \right] \right\}$$

This shows that the energy U of the electromagnetic field remains acting on the imaginary particle. In practice, this means that electromagnetic fields act on imaginary particles. Therefore, electromagnetic field of a GCC remains acting on the particles inside the GCC even when their gravitational masses reach the gravitational mass ranging between  $+0.159M_i$  to  $-0.159M_i$  and they become imaginary particles. This is very important because it means that the GCCs of a gravitational spacecraft keep when the spacecraft working becomes imaginary.

Under these conditions, the gravity accelerations on the *imaginary* spacecraft particle (due to the rest of the imaginary Universe) are given by

$$g'_{i} = \chi g_{i}$$
  $j = 1, 2, 3, ..., n$ .

Where  $\chi=M_{g\,(imaginary\,\,)}/M_{i\,(imaginary\,\,)}$  and  $g_j=-Gm_{gj\,(imaginary)}/r_j^2$ . Thus, the gravitational forces acting on the spacecraft are given by

$$F_{gj} = M_{g(imaginary)}g'_{j} =$$

$$= M_{g(imaginary)} \left(-\chi G m_{gj(imaginary)}/r_{j}^{2}\right) =$$

$$= M_{g}i\left(-\chi G m_{gj}i/r_{j}^{2}\right) = +\chi G M_{g}m_{gj}/r_{j}^{2}.$$

Note that these forces are *real*. Remind that, the Mach's principle says that the *inertial effects* upon a particle are consequence of the gravitational interaction of the particle with the rest of the Universe. Then we can conclude that the *inertial forces* upon an *imaginary* spacecraft are also *real*. Consequently, it can travel in the imaginary space-time using its thrusters.

It was shown that, *imaginary* particles can have *infinite* speed in the *imaginary* space-time [1]. Therefore, this is also the speed upper limit for the spacecraft in the imaginary space-time.

Since the gravitational spacecraft can use its thrusters after to becoming an imaginary body, then if the thrusters produce a total thrust F = 1000kN and the gravitational mass of the spacecraft is reduced from  $M_g = M_i = 10^5 kg$  down to  $M_{g} \cong 10^{-6} kg$ , the acceleration of the spacecraft will be,  $a = F/M_e \cong 10^2 ms^{-2}$ . With this acceleration the spacecraft crosses the "visible" ( $diameter = d \approx 10^{26} m$ ) in a time interval  $\Delta t = \sqrt{2d/a} \cong 1.4 \times 10^7 \, \text{m.s}^{-1} \cong 5.5 \, \text{months}$ Since the inertial effects upon the spacecraft are reduced  $\stackrel{\cdot}{M}_{\mathrm{g}}/M_{i}\cong10^{-11}$  then, in spite of the effective spacecraft acceleration  $a = 10^{12} \, m.s^{-1}$ , the effects for the crew and for the spacecraft will be equivalent to an acceleration a' given by

$$a' = \frac{M_g}{M_i} a \approx 10 m.s^{-1}$$

This is the order of magnitude of the acceleration upon of a commercial jet aircraft.

On the other hand, the travel in the *imaginary* space-time can be very safe, because there won't any material body along the trajectory of the spacecraft.

Now consider the GCCs presented in Fig. 8 (a). Note that below and above the air are the bottom and the top of the chamber. Therefore the choice of the material of the chamber is highly relevant. If the chamber is made of steel, for example, and the gravity acceleration below the chamber is g then at the bottom of the chamber, the gravity becomes  $g' = \chi_{steel} g$ ; in the air, the gravity is  $g'' = \chi_{air}g' = \chi_{air}\chi_{steel}g$ . At the top of the chamber,  $g''' = \chi_{steel} g'' = (\chi_{steel})^2 \chi_{air} g$ . Thus, out of the chamber (close to the top) the gravity acceleration becomes g". (See Fig. 8 (a)). However, for the steel at B < 300T and  $f = 1 \times 10^{-6} Hz$ , we

$$\chi_{steel} = \frac{m_{g(steel)}}{m_{i(steel)}} = \left\{1 - 2\left[\sqrt{1 + \frac{\sigma_{(steel)}B^4}{4\pi f \mu \rho_{(steel)}^2 c^2}} - 1\right]\right\} \cong 1$$

Since  $\rho_{steel} = 1.1 \times 10^6 \, S.m^{-1}$ ,  $\mu_r = 300$  and  $\rho_{(steel)} = 7800 k.m^{-3}$ .

Thus, due to  $\chi_{\mathit{steel}} \cong 1$  it follows that

$$g''' \cong g'' = \chi_{air} g' \cong \chi_{air} g$$

If instead of one GCC we have three GCC, all with steel box (Fig. 8(b)), then the gravity acceleration above the second GCC,  $g_2$  will be given by

$$g_2 \cong \chi_{air} g_1 \cong \chi_{air} \chi_{air} g$$

and the gravity acceleration above the *third* GCC,  $g_3$  will be expressed by

$$g_3 \cong \chi_{air} g'' \cong \chi_{air}^3 g$$

## III. CONSEQUENCES

These results point the possibility to convert gravitational energy rotational mechanical energy. Consider for example the system presented in Fig. 9. Basically it is a motor with massive iron rotor and a box filled with gas or plasma at ultra-low pressure (Gravity Control Cell-GCC) as shown in Fig. 9. The GCC is placed below the rotor in order to become *negative* the acceleration of gravity inside *half* of the rotor  $(g' = (\chi_{steel})^2 \chi_{air} g \cong \chi_{air} g = -ng)$ . Obviously this causes a torque T = (-F' + F)r and the rotor spins with angular velocity  $\omega$ . The average power, P, of the motor is given by

$$P = T\omega = \left[ \left( -F' + F \right) r \right] \omega \tag{36}$$

Where

$$F' = \frac{1}{2} m_g g' \qquad F = \frac{1}{2} m_g g$$

and  $m_g \cong m_i$  ( mass of the rotor ). Thus, Eq. (36) gives

$$P = (n+1)\frac{m_i g\omega r}{2} \tag{37}$$

On the other hand, we have that

$$-g' + g = \omega^2 r \tag{38}$$

Therefore the angular speed of the rotor is given by

$$\omega = \sqrt{\frac{(n+1)g}{r}} \tag{39}$$

By substituting (39) into (37) we obtain the expression of the average power of the *gravitational motor*, i.e.,

$$P = \frac{1}{2} m_i \sqrt{(n+1)^3 g^3 r}$$
 (40)

Now consider an electric generator coupling to the gravitational motor in order to produce electric energy.

Since  $\omega = 2\pi f$  then for f = 60Hzwe have  $\omega = 120 \pi rad .s^{-1} = 3600 rpm$ .

Therefore for  $\omega = 120 \pi rad.s^{-1}$  and  $n = 788 \ (B \cong 0.22T)$  the Eq. (40) tell us that we must have

$$r = \frac{(n+1)g}{\omega^2} = 0.0545m$$

Since r=R/3 and  $m_i=\rho\pi R^2h$  where  $\rho$ , R and h are respectively the mass density, the radius and the height of the rotor then for h=0.5m and  $\rho=7800Kg.m^{-3}$  (iron) we obtain

$$m_i = 327.05kg$$

Then Eq. (40) gives

$$P \cong 2.19 \times 10^5 \, watts \cong 219 \, KW \cong 294HP \tag{41}$$

This shows that the *gravitational motor* can be used to yield electric energy at large scale.

The possibility of gravity control leads to a new concept of spacecraft which is presented in Fig. 10. Due to the *Meissner effect*, the magnetic field *B* is expelled from the *superconducting shell*. The Eq. (35) shows that a magnetic field, *B*, through the *aluminum shell* of the spacecraft reduces its gravitational mass according to the following expression:

$$m_{g(Al)} = \left\{ 1 - 2 \left[ \sqrt{1 + \left( \frac{B^2}{\mu c^2 \rho_{(Al)}} n_{r(Al)} \right)^2} - 1 \right] \right\} m_{i(Al)} \quad (42)$$

If the frequency of the magnetic field is  $f=10^{-4}\,Hz$  then we have that  $\sigma_{(Al)}>>\omega\varepsilon$  since the electric conductivity of the aluminum is  $\sigma_{(Al)}=3.82\times10^7\,S.m^{-1}$ . In this case, the Eq. (11) tell us that

$$n_{r(Al)} = \sqrt{\frac{\mu c^2 \sigma_{(Al)}}{4\pi f}} \tag{43}$$

Substitution of (43) into (42) yields

$$m_{g(Al)} = \left\{ 1 - 2 \left[ \sqrt{1 + \frac{\sigma_{(Al)} B^4}{4 \pi f \mu \rho_{(Al)}^2 c^2}} - 1 \right] \right\} m_{i(Al)} \quad (44)$$

Since the mass density of the Aluminum is  $\rho_{(Al)} = 2700 \, kg \, .m^{-3}$  then the Eq. (44) can be rewritten in the following form:

$$\chi_{Al} = \frac{m_{g(Al)}}{m_{l(Al)}} = \left\{ 1 - 2 \left[ \sqrt{1 + 3.68 \times 10^{-8} B^4} - 1 \right] \right\} (45)$$

In practice it is possible to adjust B in order to become, for example,  $\chi_{Al}\cong 10^{-9}$ . This occurs to  $B\cong 76.3T$ .

(Novel superconducting magnets are able to produce up to 14.7T [10, 11]).

Then the gravity acceleration in any direction *inside* the spacecraft,  $g_l'$ , will be reduced and given by

$$g'_{l} = \frac{m_{g(Al)}}{m_{i(Al)}} g_{l} = \chi_{Al} g_{l} \cong -10^{-9} g_{l} \quad l = 1,2,...,n$$

Where  $g_l$  is the *external* gravity in the direction l. We thus conclude that the gravity acceleration inside the spacecraft becomes negligible if  $g_l << 10^9 \, m.s^{-2}$ . This means that the aluminum shell, under these conditions, works like a gravity shielding.

Consequently, the gravitational forces between anyone point inside the spacecraft with gravitational mass,  $m_{gj}$ , and another external to the spacecraft (gravitational mass  $m_{gk}$ ) are given by

$$\vec{F}_j = -\vec{F}_k = -G \frac{m_{gj} m_{gk}}{r_{jk}^2} \hat{\mu}$$

where  $m_{gk}\cong m_{ik}$  and  $m_{gj}=\chi_{Al}m_{ij}$ . Therefore we can rewrite equation above in the following form

$$\vec{F}_j = -\vec{F}_k = -\chi_{Al}G\frac{m_{ij}m_{ik}}{r_{jk}^2}\hat{\mu}$$

Note that when B = 0 the *initial* gravitational forces are

$$\vec{F}_j = -\vec{F}_k = -G \frac{m_{ij} m_{ik}}{r_{ik}^2} \hat{\mu}$$

Thus, if  $\chi_{Al} \cong -10^{-9}$  then the initial gravitational forces are reduced from  $10^9$  times and become repulsives.

According to the new expression for the *inertial forces* [1],  $\vec{F} = m_g \vec{a}$ , we see that these forces have origin in the *gravitational interaction* between a particle and the others of the Universe, just as *Mach's principle* predicts. Hence mentioned expression incorporates the Mach's principle into Gravitation Theory,

and furthermore reveals that the inertial effects upon a body can be strongly reduced by means of the decreasing of its gravitational mass.

Consequently, we conclude that if the *gravitational forces* upon the spacecraft are reduced from  $10^9$  times then also the *inertial forces* upon the spacecraft will be reduced from  $10^9$  times when  $\chi_{Al} \cong -10^{-9}$ . Under these conditions, the inertial effects on the crew would be strongly decreased. Obviously this leads to a new concept of aerospace flight.

Inside the spacecraft the gravitational forces between the dielectric with gravitational mass,  $M_g$  and the man (gravitational mass,  $m_g$ ), when B=0 are

$$\vec{F}_{m} = -\vec{F}_{M} = -G \frac{M_{g} m_{g}}{r^{2}} \hat{\mu}$$
 (46)

or

$$\vec{F}_{m} = -G \frac{M_{g}}{r^{2}} m_{g} \hat{\mu} = -m_{g} g_{M} \hat{\mu}$$
 (47)

$$\vec{F}_{M} = +G \frac{m_{g}}{r^{2}} M_{g} \hat{\mu} = +M_{g} g_{m} \hat{\mu}$$
 (48)

If the *superconducting box* under  $M_g$  (Fig. 10) is filled with *air* at ultra-low pressure (3×10<sup>-12</sup> torr, 300K for example) then, when  $B \neq 0$ , the gravitational mass of the *air* will be reduced according to (35). Consequently, we have

$$g'_{M} = (\chi_{steel})^{2} \chi_{air} g_{M} \cong \chi_{air} g_{M}$$
 (49)

$$g'_{m} = (\chi_{steel})^{2} \chi_{air} g_{m} \cong \chi_{air} g_{m}$$
 (50)

Then the forces  $\vec{F}_m$  and  $\vec{F}_M$  become

$$\vec{F}_m = -m_g \left( \chi_{air} g_M \right) \hat{\mu} \tag{51}$$

$$\vec{F}_{M} = +M_{g} \left( \chi_{air} g_{m} \right) \hat{\mu} \tag{52}$$

Therefore if  $\chi_{air} = -n$  we will have

$$\vec{F}_m = +nm_g g_M \hat{\mu} \tag{53}$$

$$\vec{F}_{M} = -nM_{g}g_{m}\hat{\mu} \tag{54}$$

Thus,  $\vec{F}_m$  and  $\vec{F}_M$  become *repulsive*. Consequently, the man inside the

spacecraft is subjected to a gravity acceleration given by

$$\vec{a}_{man} = ng_M \hat{\mu} = -\chi_{air} G \frac{M_g}{r^2} \hat{\mu}$$
 (55)

Inside the GCC we have,

$$\chi_{air} = \frac{m_{g(air)}}{m_{i(air)}} = \left\{ 1 - 2 \left[ \sqrt{1 + \frac{\sigma_{(air)}B^4}{4\pi f \mu \rho_{(air)}^2 c^2}} - 1 \right] \right\}$$
 (56)

By ionizing the air inside the GCC (Fig. 10), for example, by means of a radioactive material, it is possible to increase the air conductivity inside the GCC up to  $\sigma_{(air)} \cong 10^6 \, \text{S.m}^{-1}$ . Then for  $f = 10 \, Hz$ ;  $\rho_{(air)} = 4.94 \times 10^{-15} \, kg.m^{-3}$  (Air at  $3 \times 10^{-12}$  torr, 300K) and we obtain

$$\chi_{air} = \left\{ 2 \left[ \sqrt{1 + 2.8 \times 10^{21} B^4} - 1 \right] - 1 \right\}$$
(57)

For  $B = B_{GCC} = 0.1T$  (note that, due to the *Meissner effect*, the magnetic field  $B_{GCC}$  stay confined inside the superconducting box) the Eq. (57) yields

$$\chi_{air} \cong -10^9$$

Since there is no magnetic field through the *dielectric* presented in Fig.10 then,  $M_g \cong M_i$ . Therefore if  $M_g \cong M_i = 100 \mathrm{Kg}$  and  $r = r_0 \cong 1m$  the gravity acceleration upon the man, according to Eq. (55), is

$$a_{man} \cong 10 m.s^{-1}$$

Consequently it is easy to see that this system is ideal to yield artificial gravity inside the spacecraft in the case of *interstellar travel*, when the gravity acceleration out of the spacecraft - due to the Universe - becomes negligible.

The *vertical* displacement of the spacecraft can be produced by means of *Gravitational Thrusters*. A schematic diagram of a Gravitational Thruster is shown in Fig.11. The Gravitational

Thrusters can also provide the *horizontal* displacement of the spacecraft.

The concept of Gravitational Thruster results from the theory of the *Gravity Control Battery*, showed in Fig. 8 (b). Note that the number of GCC increases the thrust of the thruster. For example, if the thruster has *three* GCCs then the gravity acceleration upon the gas sprayed inside the thruster will be *repulsive* in respect to  $M_g$  (See Fig. 11) and given by

$$a_{gas} = (\chi_{air})^3 (\chi_{steel})^4 g \cong -(\chi_{air})^3 G \frac{M_g}{r_0^2}$$

Thus, if inside the GCCs,  $\chi_{air} \cong -10^9$  (See Eq. 56 and 57) then the equation above gives

$$a_{gas} \cong +10^{27} G \frac{M_i}{r_0^2}$$

For  $M_i \cong 10 kg$ ,  $r_0 \cong 1 m$  and  $m_{gas} \cong 10^{-12} kg$  the thrust is

$$F = m_{gas} a_{gas} \cong 10^5 N$$

Thus, the Gravitational Thrusters are able to produce strong thrusts.

It is important to note that if F is force produced by a thruster then the spacecraft acquires acceleration  $a_{spacecraft}$  given by [1]

$$a_{spacecraft} = \frac{F}{M_{g(spacecraft)}} = \frac{F}{\chi_{Al} M_{i(inside)} + m_{i(Al)}}$$

Therefore if  $\chi_{Al} \cong 10^{-9}$ ;  $M_{i(inside)} = 10^4 Kg$  and  $m_{i(Al)} = 100 Kg$  (inertial mass of the aluminum shell) then it will be necessary F = 10kN to produce

$$a_{spacecraft} = 100 m.s^{-2}$$

Note that the concept of Gravitational Thrusters leads directly to the *Gravitational Turbo Motor* concept (See Fig. 12).

Let us now calculate the gravitational forces between two very close *thin* layers of the *air* around the spacecraft. (See Fig. 13).

The gravitational force  $dF_{12}$  that  $dm_{g1}$  exerts upon  $dm_{g2}$ , and the gravitational force  $dF_{21}$  that  $dm_{g2}$  exerts upon  $dm_{g1}$  are given by

$$d\vec{F}_{12} = d\vec{F}_{21} = -G \frac{dm_{g2}dm_{g1}}{r^2} \hat{\mu}$$
 (58)

Thus, the gravitational forces between the *air layer* 1, gravitational mass  $m_{g1}$ , and the *air layer* 2, gravitational mass  $m_{g2}$ , around the spacecraft are

$$\vec{F}_{12} = -\vec{F}_{21} = -\frac{G}{r^2} \int_0^{m_{g1}} \int_0^{m_{g2}} dm_{g1} dm_{g2} \hat{\mu} =$$

$$= -G \frac{m_{g1} m_{g2}}{r^2} \hat{\mu} = -\chi_{air} \chi_{air} G \frac{m_{i1} m_{i2}}{r^2} \hat{\mu}$$
(59)

At 100km altitude the air pressure is  $5.69 \times 10^{-3} torr$  and  $\rho_{(air)} = 5.99 \times 10^{-6} kgm^{-3}$  [12]. By ionizing the air surround the spacecraft, for example, by means of an oscillating electric field,  $E_{osc}$ , starting from the surface of the spacecraft ( See Fig. 13) it is possible to increase the *air conductivity* near the spacecraft up to  $\sigma_{(air)} \cong 10^6 S.m^{-1}$ . Since f = 1Hz and, in this case  $\sigma_{(air)} >> \omega \varepsilon$ , then, according to

Eq. (11),  $n_r = \sqrt{\mu \sigma_{(air)} c^2 / 4\pi f}$  . From Eq.(56) we thus obtain

$$\chi_{air} = \frac{m_{g(air)}}{m_{f(air)}} = \left\{ 1 - 2 \sqrt{1 + \frac{\sigma_{(air)}B^4}{4\pi f \mu_0 \rho_{(air)}^2 c^2}} - 1 \right\}$$
 (60)

Then for B = 763T the Eq. (60) gives

$$\chi_{air} = \left\{1 - 2\left[\sqrt{1 + \sim 10^4 B^4} - 1\right]\right\} \cong -10^8$$
 (61)  
By substitution of  $\chi_{air} \cong -10^8$  into Eq., (59) we get

$$\vec{F}_{12} = -\vec{F}_{21} = -10^{16} G \frac{m_{i1} m_{i2}}{r^2} \hat{\mu}$$
 (62)

If  $m_{i1} \cong m_{i2} = \rho_{air} V_1 \cong \rho_{air} V_2 \cong 10^{-8} kg$ , and  $r = 10^{-3} m$  we obtain

$$\vec{F}_{12} = -\vec{F}_{21} \cong -10^{-4} N \tag{63}$$

These forces are much more intense than the *inter-atomic forces* (the forces which maintain joined atoms, and molecules that make the solids and liquids) whose intensities, according to the Coulomb's law, is of the order of 1-1000×10<sup>-8</sup>N.

Consequently, the air around the spacecraft will be strongly compressed upon their surface, making an "air shell" that will accompany the spacecraft during its displacement and will protect the aluminum shell of the direct attrition with the Earth's atmosphere.

In this way, during the flight, the attrition would occur just between the "air shell" and the atmospheric air around her. Thus, the spacecraft would stay free of the thermal effects that would be produced by the direct attrition of the aluminum shell with the Earth's atmosphere.

Another interesting effect produced by the magnetic field B of the spacecraft is the possibility of to lift a body from the surface of the Earth to the spacecraft as shown in Fig. 14. By ionizing the air surround the spacecraft, by means of an oscillating electric field,  $E_{osc}$ , the air conductivity near the spacecraft can reach, for example,  $\sigma_{(air)} \cong 10^6 \, S.m^{-1}$ . Then for f = 1Hz; B = 40.8T and  $\rho_{(air)} \cong 1.2kg.m^{-3}$  (300K and 1 atm) the Eq. (56) yields

$$\chi_{air} = \left\{1 - 2\left[\sqrt{1 + 4.9 \times 10^{-7} B^4} - 1\right]\right\} \cong -0.1$$

Thus, the weight of the body becomes

 $P_{body} = m_{g(body)}g = \chi_{air}m_{i(body)}g = m_{i(body)}g'$ Consequently, the body will be lifted on the direction of the spacecraft with acceleration

$$g' = \chi_{air} g \cong +0.98 \text{m.s}^{-1}$$

Let us now consider an important

aspect of the flight dynamics of a Gravitational Spacecraft.

Before starting the flight, the gravitational mass of the spacecraft,  $M_{g}$ , must be strongly reduced, by means of a gravity control system, in order to produce — with a weak thrust  $\vec{F}$ , a strong acceleration,  $\vec{a}$ , given by [1]

$$\vec{a} = \frac{\vec{F}}{M_g}$$

In this way, the spacecraft could be strongly accelerated and quickly to reach very high speeds near speed of light.

If the gravity control system of the spacecraft is *suddenly* turned off, the *gravitational mass* of the spacecraft becomes immediately equal to its *inertial mass*,  $M_i$ ,  $\left(M_g' = M_i\right)$  and the velocity

 $\vec{V}$  becomes equal to  $\vec{V}'.$  According to the  $\it Momentum$  Conservation Principle, we have that

$$M_{g}V = M'_{g}V'$$

Supposing that the spacecraft was traveling in space with speed  $V \approx c$ , and that its gravitational mass it was  $M_g = 1Kg$  and  $M_i = 10^4\,Kg$  then the velocity of the spacecraft is reduced to

$$V' = \frac{M_g}{M'_g}V = \frac{M_g}{M_i}V \approx 10^{-4}c$$

Initially, when the velocity of the spacecraft is  $\vec{V}$ , its kinetic energy is  $E_{\!_{k}} = \!\! \left(\! M_{\!_{g}} \! - \! m_{\!_{g}} \! \right) \!\! c^2$ . Where  $M_{\!_{g}} = m_{\!_{g}} / \sqrt{1 - V^2/c^2}$ . At the instant in which the gravity control system of the spacecraft is turned off, the kinetic energy becomes  $E_{\!_{k}}' = \! \left(\! M_{\!_{g}}' \! - \! m_{\!_{g}}' \! \right) \!\! c^2$ . Where  $M_{\!_{g}}' = m_{\!_{g}}' / \sqrt{1 - V'^2/c^2}$ .

We can rewritten the expressions of  $E_k$  and  $E_k'$  in the following form

$$E_k=ig(M_gV-m_gVig)rac{c^2}{V}$$
 
$$E_k'=ig(M_g'V'-m_g'V'ig)rac{c^2}{V'}$$
 Substitution of  $M_gV=M_g'V'=p$  ,

 $m_{\!g}V = p\sqrt{1-V^2/c^2}$  and  $m_{\!g}'V' = p\sqrt{1-V'^2/c^2}$  into the equations of  $E_k$  and  $E_k'$  gives

$$E_k = \left(1 - \sqrt{1 - V^2/c^2}\right) \frac{pc^2}{V}$$

$$E'_{k} = \left(1 - \sqrt{1 - V'^{2}/c^{2}}\right) \frac{pc^{2}}{V'}$$

Since  $V \approx c$  then follows that

$$E_{\nu} \approx pc$$

On the other hand, since  $V' \ll c$  we get  $E'_{k} = \left(1 - \sqrt{1 - {V'}^{2}/c^{2}}\right) \frac{pc^{2}}{V'} =$ 

$$\cong \left(1 - \frac{1}{1 + \frac{V'^2}{2c^2} + \dots}\right) \frac{pc^2}{V'} \cong \left(\frac{V'}{2c}\right) pc$$

Therefore we conclude that  $E_k >> E_k'$ . Consequently, when the gravity control system of the spacecraft is turned off, occurs an *abrupt* decrease in the kinetic energy of the spacecraft,  $\Delta E_k$ , given by

$$\Delta E_k = E_k - E_k' \approx pc \approx M_g c^2 \approx 10^{17} J$$

By comparing the energy  $\Delta E_k$  with the *inertial energy* of the spacecraft,  $E_i = M_i c^2$ , we conclude that

$$\Delta E_k \approx \frac{M_g}{M_i} E_i \approx 10^{-4} M_i c^2$$

The energy  $\Delta E_k$  (several *megatons*) must be released in very short time interval. It is approximately the same amount of energy that would be released in the case of collision of the spacecraft<sup>‡</sup>. However, the situation is very different of a collision  $(M_g)$  just becomes suddenly equal to  $M_i$ ), and possibly the energy  $\Delta E_k$  is converted into a *High Power Electromagnetic Pulse*.

Obviously this electromagnetic pulse (EMP) will induce heavy currents in all electronic equipment that mainly contains semiconducting and conducting materials. This produces immense heat that melts the circuitry inside. As such, while not being directly responsible for the loss of lives, these EMP are capable of disabling electric/electronic systems. Therefore, we possibly have a new type electromagnetic bomb. electromagnetic bomb or E-bomb is a well-known weapon designed to disable electric/electronic systems on a wide scale with an intense electromagnetic pulse.

Based on the theory of the GCC it is also possible to build a *Gravitational Press* of *ultra-high* pressure as shown in Fig.15.

The chamber 1 and 2 are GCCs with air at  $1\times 10^{-4} \text{torr}$ , 300K  $\left(\sigma_{(air)}\approx 10^6\,S.m^{-1}; \rho_{(air)}=5\times 10^{-8}\,kg.m^{-3}\right)$ . Thus, for f=10Hz and B=0.107T we have

$$\chi_{air} = \left\{ 1 - 2 \left[ \sqrt{1 + \frac{\sigma_{(air)} B^4}{4\pi f \mu_0 \rho_{(air)}^2 c^2}} - 1 \right] \right\} \cong -118$$

The gravity acceleration above the air of the chamber 1 is

$$\vec{g}_1 = \chi_{stell} \chi_{air} g \hat{\mu} \cong +1.15 \times 10^3 \hat{\mu}$$
 (64) Since, in this case,  $\chi_{steel} \cong 1$ ;  $\hat{\mu}$  is an *unitary vector* in the opposite direction of  $\vec{g}$ .

Above the air of the chamber 2 the gravity acceleration becomes

$$\vec{g}_2 = (\chi_{stell})^2 (\chi_{air})^2 g \hat{\mu} \cong -1.4 \times 10^5 \hat{\mu}$$
 (65)

Therefore the *resultant* force  $\vec{R}$  acting on  $m_2$ ,  $m_1$  and m is

<sup>&</sup>lt;sup>‡</sup> In this case, the collision of the spacecraft would release  $\approx 10^{17} \text{J}$  (several megatons) and it would be similar to a powerful *kinetic weapon*.

$$\vec{R} = \vec{F}_2 + \vec{F}_1 + \vec{F} = m_2 \vec{g}_2 + m_1 \vec{g}_1 + m \vec{g} =$$

$$= -1.4 \times 10^5 m_2 \hat{\mu} + 1.15 \times 10^3 m_1 \hat{\mu} - 9.81 m \hat{\mu} =$$

$$\approx -1.4 \times 10^5 m_2 \hat{\mu}$$
(66)

where

$$m_2 = \rho_{steel} V_{disk2} = \rho_{steel} \left( \frac{\pi}{4} \phi_{inn}^2 H \right)$$
 (67)

Thus, for  $\rho_{steel} \cong 10^4 \, kg \, m^{-3}$  we can write that

$$F_2 \cong 10^9 \phi_{inn}^2 H$$

For the steel  $\tau \cong 10^5 \, kg.cm^{-2} = 10^9 \, kg.m^{-2}$  consequently we must have  $F_2/S_\tau < 10^9 \, kg.m^{-2}$  ( $S_\tau = \pi \phi_{inn} H$  see Fig.15). This means that

$$\frac{10^9 \phi_{inn}^2 H}{\pi \phi_{inn} H} < 10^9 \, kg \, .m^{-2}$$

Then we conclude that

$$\phi_{inn} < 3.1m$$

For  $\phi_{inn} = 2m$  and H = 1m the Eq. (67) gives

$$m_2 \cong 3 \times 10^4 kg$$

Therefore from the Eq. (66) we obtain  $R \cong 10^{10} N$ 

Consequently, in the area  $S = 10^{-4} m^2$  of the Gravitational Press, the pressure is

$$p = \frac{R}{S} \cong 10^{14} \, N.m^{-2}$$

This enormous pressure is much greater than the pressure in the center of the Earth  $(3.617 \times 10^{11} N.m^{-2})$  [13]. It is near of the gas pressure in the center of the sun  $(2 \times 10^{16} N.m^{-2})$ . Under the action of such intensities new states of matter created and astrophysical phenomena may be simulated in the lab for the first time, e.g. supernova explosions. Controlled thermonuclear fusion by inertial confinement, fast nuclear ignition for energy gain, novel collective acceleration schemes particles and the numerous variants of material processing constitute examples of progressive applications of such Press Gravitational of ultra-high pressure.

The GCCs can also be applied on generation and detection of *Gravitational Radiation*.

Consider a cylindrical GCC (GCC antenna) as shown in Fig.16 (a). The gravitational mass of the air inside the GCC is

$$m_{g(air)} = \left\{ 1 - 2 \left[ \sqrt{1 + \frac{\sigma_{(air)}B^4}{4\pi f \mu \rho_{(air)}^2 c^2}} - 1 \right] \right\} m_{i(air)} \quad (68)$$

By varying B one can varies  $m_{g(air)}$  and consequently to vary the gravitational field generated by  $m_{g(air)}$ , producing then gravitational radiation. Then a GCC can work like a *Gravitational Antenna*.

Apparently, Newton's theory of gravity had no gravitational waves because, if a gravitational field changed in some way, that change took place instantaneously everywhere in space, and one can think that there is not a wave in this case. However, we have already seen that the gravitational interaction can be repulsive, besides attractive. Thus, as with electromagnetic interaction, the gravitational interaction must be produced by the exchange of "virtual" quanta of spin 1 and mass null, i.e., the gravitational "virtual" quanta (graviphoton) must have spin 1 and not Consequently, the fact of a change in gravitational field instantaneously everywhere in space occurs simply due to the speed of the graviphoton to be infinite. It is known that there is no speed limit for "virtual" contrary, photons. On the electromagnetic ("virtual" quanta photons) could not communicate the electromagnetic interaction an infinite distance.

Thus, there are *two types* of gravitational radiation: the *real* and *virtual*, which is constituted of graviphotons; the *real* gravitational waves are ripples in the space-time generated by *gravitational field* changes. According to Einstein's theory of gravity the velocity of propagation of these waves is equal to the speed of light (*c*).

Unlike the electromagnetic waves the real gravitational waves have low interaction with matter and consequently low scattering. Therefore gravitational waves are suitable as a of transmitting information. However, when the distance between transmitter and receiver is too large. for example of the order of magnitude of several light-years, the transmission of information by means of gravitational waves becomes impracticable due to the long time necessary to receive the information. On the other hand, there is no delay during the transmissions by means of *virtual* gravitational radiation. In addition the scattering of this radiation is null. Therefore the virtual gravitational radiation is very suitable as a means of transmitting information at any distances including astronomical distances.

As concerns detection of the virtual gravitational radiation from GCC antenna, there are many options. Due to Resonance Principle a similar GCC antenna (receiver) tuned at the same frequency can absorb energy from an incident virtual gravitational radiation (See Fig.16 (b)). Consequently, the gravitational mass of the air inside the GCC receiver will vary such as the gravitational mass of the air inside the GCC transmitter. This will induce a magnetic field similar to the magnetic field of the GCC transmitter and therefore the current through the coil inside the GCC receiver will have the same characteristics of the current through the coil inside the GCC transmitter. However, the volume and pressure of the air inside the two GCCs must be exactly the same; also the type and the *quantity of atoms* in the air inside the two GCCs must be exactly the same. Thus, the GCC antennas are simple but they are not easy to build.

Note that a GCC antenna radiates graviphotons and gravitational waves simultaneously (Fig. 16 (a)). Thus, it is not only a gravitational antenna: it is a Quantum Gravitational Antenna because

can also emit and detect iŧ gravitational "virtual" quanta (graviphotons), which, in turn, can transmit information instantaneously from any distance in the Universe without scattering.

Due to the difficulty to build two similar GCC antennas and, considering that the electric current in the receiver antenna can be detectable even if the gravitational mass of the nuclei of the antennas are not strongly reduced, then we propose to replace the gas at the nuclei of the antennas by a thin dielectric lamina. The dielectric lamina with exactly  $10^8$  atoms  $(10^3 \text{atoms} \times 10^3 \text{atoms} \times$ 10<sup>2</sup> atoms) is placed between the plates (electrodes) as shown in Fig. 17. When the virtual gravitational radiation strikes the dielectric lamina. nogu gravitational mass varies similarly to the gravitational mass of the lamina of the transmitter antenna. inducing an electromagnetic field (E, B)similar to the transmitter antenna. Thus, the electric current in the receiver antenna will have the characteristics of the current in the transmitter antenna. In this way, it is then possible to build two similar antennas whose nuclei have the same volumes and the same types and quantities of atoms.

Note that the Quantum Gravitational Antennas can also be used to transmit electric power. It is easy to see that the Transmitter and Receiver 17(a)) can work with strong voltages and electric currents. This means that strong electric power can be transmitted among Quantum Gravitational Antennas. This obviously solves the problem of wireless electric power transmission.

The existence of *imaginary masses* has been predicted in a previous work [1]. Here we will propose a method and a device using GCCs for obtaining *images* of *imaginary bodies*.

It was shown that the *inertial* imaginary mass associated to an *electron* is given by

$$m_{ie(ima)} = \frac{2}{\sqrt{3}} \left(\frac{hf}{c^2}\right) i = \frac{2}{\sqrt{3}} m_{ie(real)} i$$

Assuming that the correlation between the gravitational mass and the inertial mass (Eq.6) is the same for both imaginary and real masses then follows that the *gravitational* imaginary mass associated to an *electron* can be written in the following form:

$$m_{ge(ima)} = \left\{1 - 2\left[\sqrt{1 + \left(\frac{U}{m_i c^2} n_r\right)^2} - 1\right]\right\} m_{ie(ima)}$$

Thus, the gravitational *imaginary* mass associated to matter can be reduced, made negative and increased, just as the gravitational real mass.

It was shown that also *photons* have imaginary mass. Therefore, the imaginary mass can be associated or *not* to the matter.

In a general way, the gravitational forces between two gravitational imaginary masses are then given by

$$\vec{F} = -\vec{F} = -G \frac{\left(iM_g\right)\left(im_g\right)}{r^2} \hat{\mu} = +G \frac{M_g m_g}{r^2} \hat{\mu}$$

Note that these forces are *real* and *repulsive*.

Now consider a gravitational imaginary mass,  $m_{g(ima)}=im_{g}$ , not associated with matter (like the gravitational imaginary mass associated to the photons) and another gravitational imaginary mass  $M_{g(ima)}=iM_{g}$  associated to a material body.

Any material body has an imaginary mass associated to it, due to the existence of imaginary masses associated to the electrons. We will choose a *quartz crystal* (for the material body with gravitational imaginary mass  $M_{g(ima)} = iM_g$ ) because quartz crystals are widely used to detect forces (piezoelectric effect).

By using GCCs as shown in Fig. 18(b) and Fig. 18(c), we can increase the

gravitational acceleration,  $\vec{a}$ , produced by the imaginary mass  $im_g$  upon the crystals. Then it becomes

$$a = -\chi_{air}^3 G \frac{m_g}{r^2}$$

As we have seen, the value of  $\chi_{air}$  can be increased up to  $\chi_{air} \cong -10^9$  (See Eq.57). Note that in this case, the gravitational forces become *attractive*. In addition, if  $m_g$  is not small, the gravitational forces between the imaginary body of mass  $im_g$  and the crystals can become sufficiently intense to be easily detectable.

Due to the piezoelectric effect, the gravitational force acting on the crystal will produce a voltage proportional to its intensity. Then consider a board with hundreds micro-crystals behind a set of GCCs, as shown in Fig.18(c). By amplifying the voltages generated in each micro-crystal and sending to an appropriated data acquisition system, it will be thus possible to obtain an image of the imaginary body of mass  $m_{g(ima)}$  placed in front of the board.

In order to decrease strongly the gravitational effects produced by bodies placed behind the imaginary body of mass  $im_g$ , one can put five GCCs making a *Gravitational Shielding* as shown in Fig.18(c). If the GCCs are filled with air at 300K and  $3\times10^{-12}torr$ . Then  $\rho_{air}=4.94\times10^{15}kgm^3$  and  $\sigma_{air}\cong1\times10^{14}Sm^4$ . Thus, for f=60Hz and  $B\cong0.7T$  the Eq. (56) gives

$$\chi_{air} = \frac{m_{g(air)}}{m_{i(air)}} = \left\{1 - 2\left[\sqrt{1 + 5B^4} - 1\right]\right\} \cong -10^{-2}$$

For  $\chi_{air} \cong 10^{-2}$  the gravitational shielding presented in Fig.18(c) will reduce any value of g to  $\chi_{air}^5 g \cong 10^{-10} g$ . This will be sufficiently to reduce strongly the gravitational effects proceeding from both sides of the gravitational shielding.

## **IV.CONCLUSION**

The process described here is clearly the better way in order to control the gravity. This is because the Gravity Control Cell in this case is very easy to be built, the cost is low and it works at ambient temperature. The Gravity Control is the starting point for the generation of and detection of Virtual Gravitational Radiation (Quantum Gravitational Transceiver) also for the construction of the Gravitational Motor and the Gravitational Spacecraft which includes the system for generation of artificial gravity presented in Fig.10 and Gravitational Thruster (Fig.11). While the Gravitational Transceiver leads to a new concept in Telecommunication, the Gravitational Motor changes the paradigm of energy conversion and the Gravitational Spacecraft points to a new concept in aerospace flight.

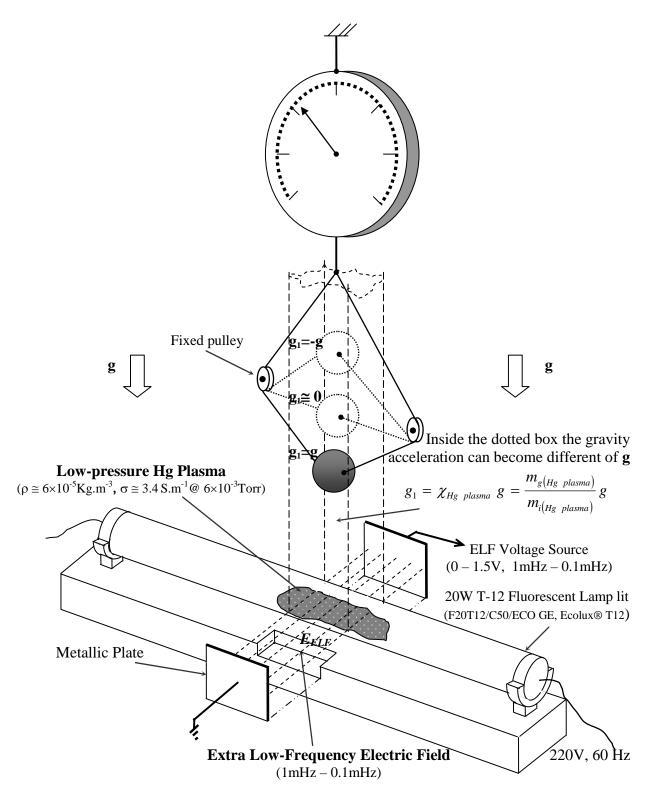


Fig. 1 – Gravitational Shielding Effect by means of an ELF electric field through low- pressure Hg Plasma.

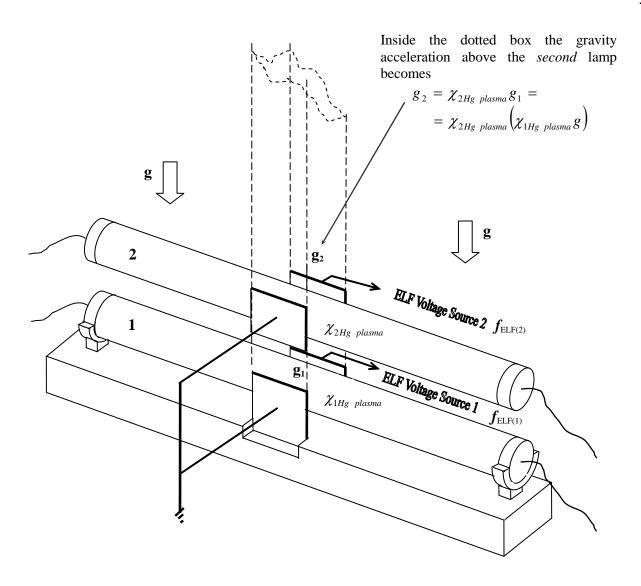


Fig. 2 – Gravity acceleration above a *second* fluorescent lamp.

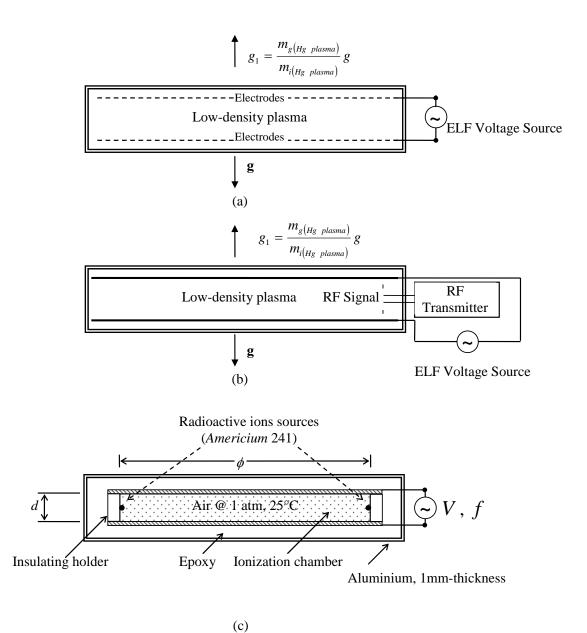
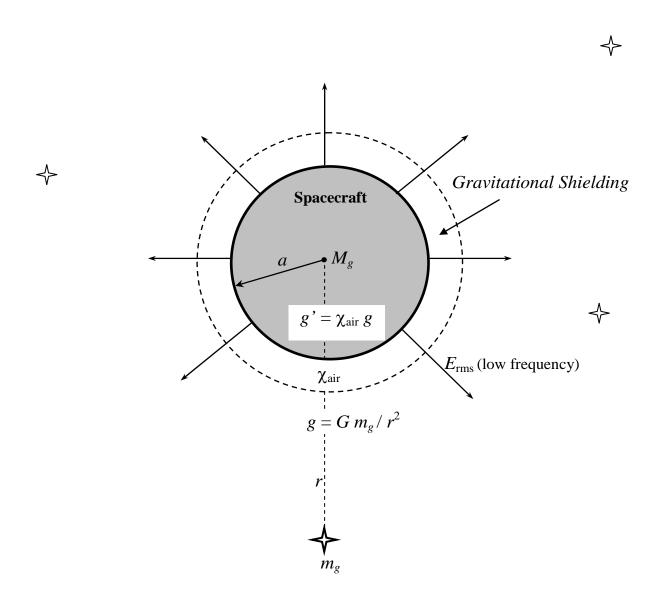


Fig. 3 – Schematic diagram of Gravity Control Cells (GCCs).

(a) GCC where the ELF electric field and the ionizing electric field can be the same. (b) GCC where the plasma is ionized by means of a RF signal. (c) GCC filled with *air* (at ambient temperature and 1 atm) strongly ionized by means of alpha particles emitted from radioactive ions sources (Am 241, *half-life* 432 years). Since the electrical conductivity of the ionized air depends on the amount of ions then it can be strongly increased by increasing the amount of Am 241 in the GCC. This GCC has 36 radioactive ions sources each one with  $1/5000^{th}$  of gram of Am 241, conveniently positioned around the ionization chamber, in order to obtain  $\sigma_{air} \cong 10^3 \, S.m^{-1}$ .



The gravity accelerations on the spacecraft (due to the rest of the Universe) can be controlled by means of the gravitational shielding, i.e.,

$$g'_{i} = \chi_{air} g_{i}$$
  $i = 1, 2, 3 ... n$ 

Thus,

$$F_{is} = F_{si} = M_g g'_i = M_g (\chi_{air} g_i)$$

 $F_{\rm is} = F_{\rm si} = M_g \ g \ {\rm 'i} = M_g \ (\chi_{\rm air} \ g_{\rm i})$  Then the inertial forces acting on the spacecraft (s) can be strongly reduced. According to the *Mach's* principle this effect can reduce the inertial properties of the spacecraft and consequently, leads to a new concept of spacecraft and aerospace flight.

Fig. 4 – Gravitational Shielding surround a Spherical Spacecraft.

$V = V_0$	t = T/4		$oldsymbol{E}_{ ext{ELF (1)}}$	$f_{\scriptscriptstyle{\mathrm{ELF}(1)}}$	$g_1/g$		$oldsymbol{E}_{ ext{ELF (2)}}$	$E_{ ext{ELF (2)}}$ $f_{ ext{ELF (2)}}$		$g_2/g$	
(Volts)	(s)	( min)	(V/m)	(mHz)	Exp.	Teo.	(V/m)	(mHz)	Exp.	Teo.	
1.0 V	250	4.17	24.81	1	-	0.993	24.81	1	-	0.986	
	312.5	5.21	24.81	0.8	-	0.986	24.81	0.8	-	0.972	
	416.6	6.94	24.81	0.6	-	0.967	24.81	0.6	-	0.935	
	625	10.42	24.81	0.4	-	0.890	24.81	0.4	-	0.792	
	1250	20.83	24.81	0.2	-	0.240	24.81	0.2	-	0.058	
1.5V	250	4.17	37.22	1	-	0.964	37.22	1	-	0.929	
	312.5	5.21	37.22	0.8	-	0.930	37.22	0.8	-	0.865	
	416.6	6.94	37.22	0.6	-	0.837	37.22	0.6	-	0.700	
	625	10.42	37.22	0.4	-	0,492	37.22	0.4	-	0.242	
	1250	20.83	37.22	0.2	-	-1,724	37.22	0.2	-	2.972	

Table 1 – Theoretical Results.

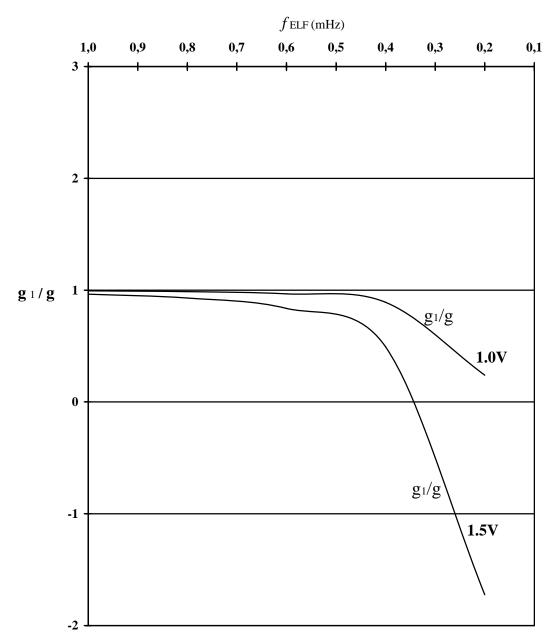


Fig. 5- Distribution of the correlation  $\,{\bf g}{\scriptscriptstyle {\rm I}}/\,{\bf g}\,$  as a function of  $f_{\rm ELF}$ 

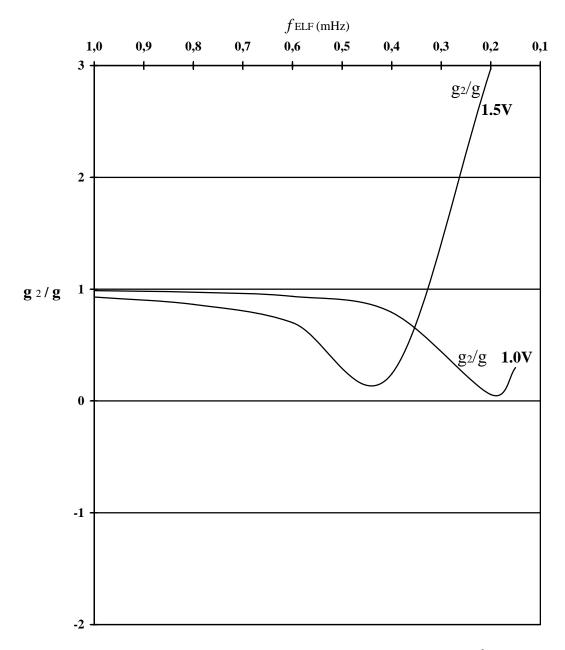


Fig. 6- Distribution of the correlation  $\,{\bf g_2}/\,{\bf g}\,$  as a function of  $f_{\rm ELF}$ 

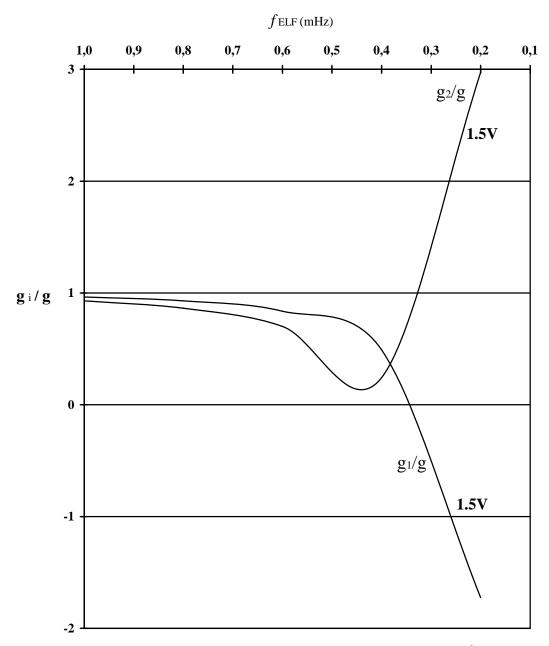
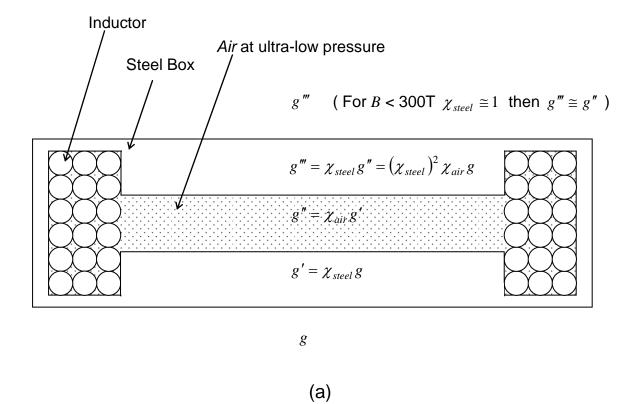


Fig. 7- Distribution of the correlations  $\,{\bf g_i}\,{\it I}\,{\bf g}\,$  as a function of  $f_{\rm ELF}$ 



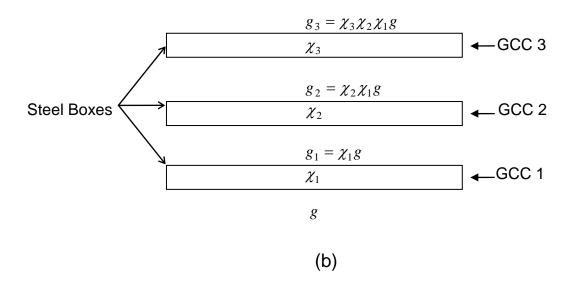
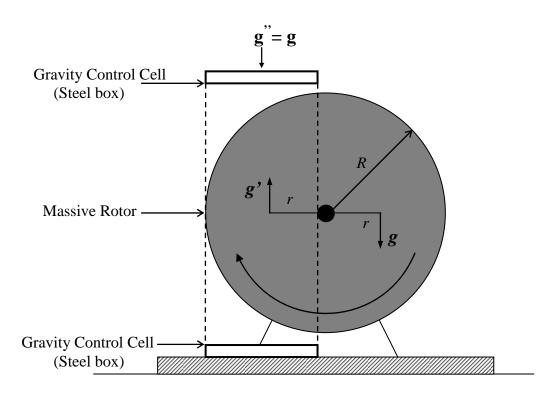


Fig. 8 – (a) Gravity Control Cell (GCC) filled with *air* at ultra-low pressure. (b) Gravity Control Battery (Note that if  $\chi_1 = \chi_2^{-1} = -1$  then g'' = g)



g

Note that  $g' = (\chi_{steel})^2 \chi_{air} g$  and  $g'' = (\chi_{steel})^4 (\chi_{air})^2 g$  therefore for  $\chi_{steel} \cong 1$  and  $\chi_{air(1)} = \chi_{air(2)}^{-1} = -n$  we get  $g' \cong -ng$  and g'' = g

Fig. 9 – The Gravitational Motor

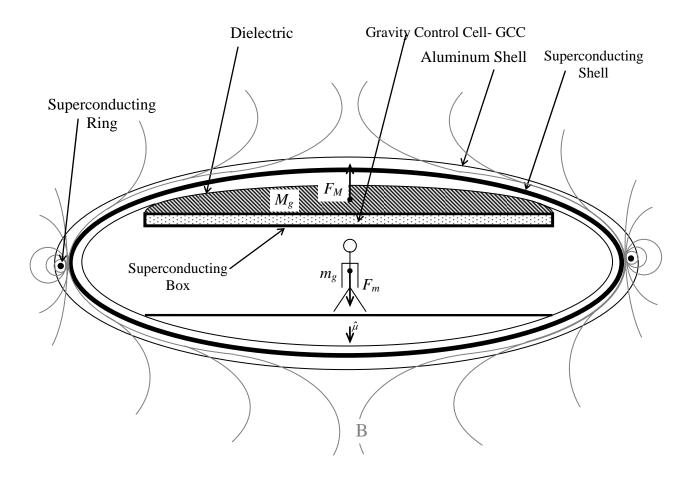
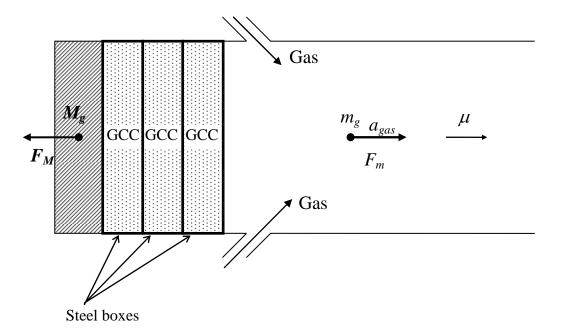


Fig. 10 – The Gravitational Spacecraft – Due to the *Meissner effect*, the magnetic field B is expelled from the *superconducting shell*. Similarly, the magnetic field  $B_{GCC}$ , of the GCC stay confined inside the *superconducting box*.



 $Fig.\ 11-The\ Gravitational\ Thruster$ 

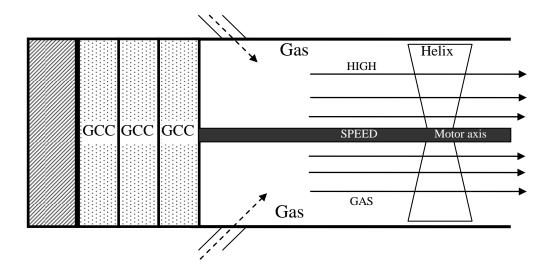


Fig. 12 - The Gravitational Turbo Motor – The gravitationally accelerated gas, by means of the GCCs, propels the helix which movies the motor axis.

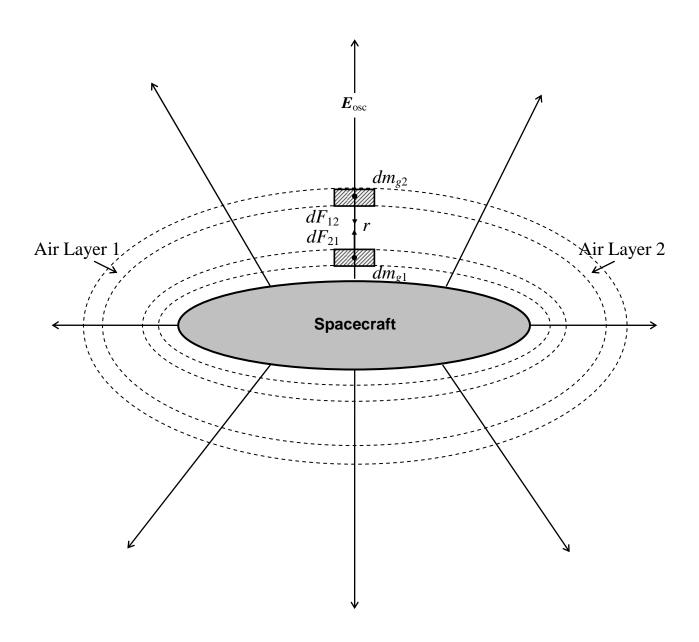


Fig. 13 – Gravitational forces between two layers of the "air shell". The electric field  $E_{\rm osc}$  provides the *ionization* of the air.

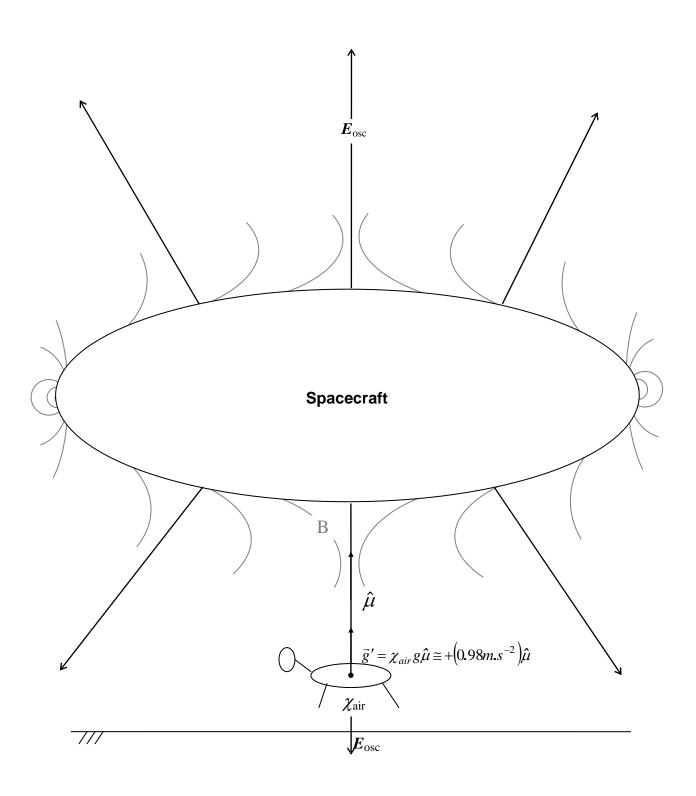


Fig. 14 – The Gravitational Lifter

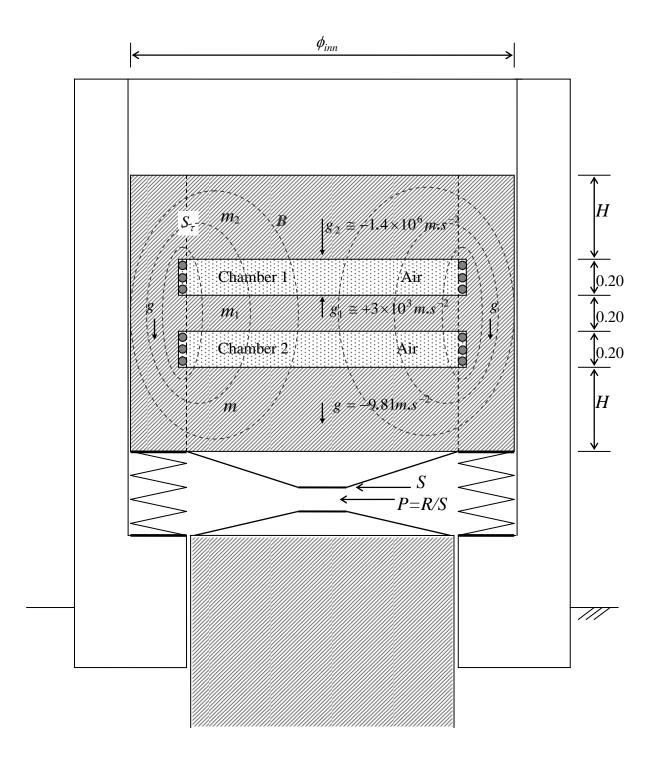


Fig. 15 – Gravitational Press

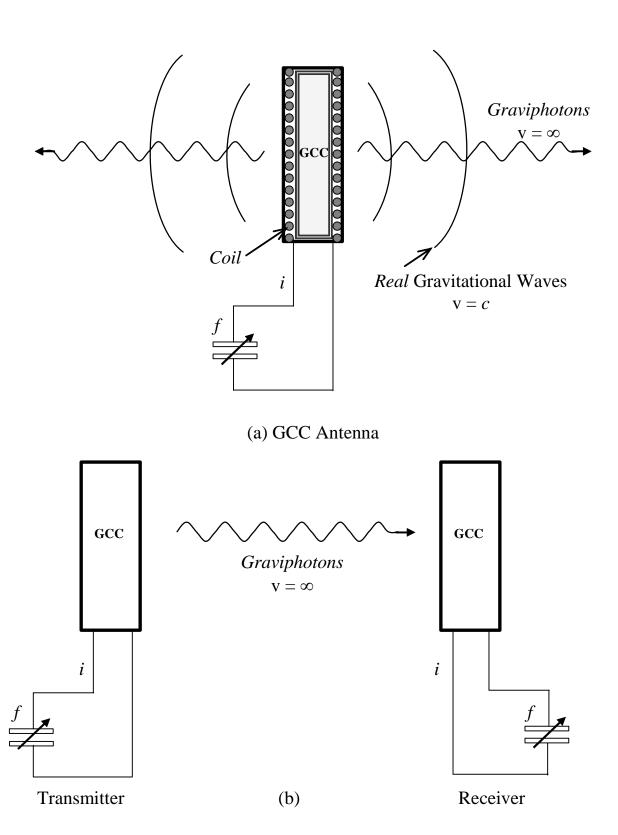
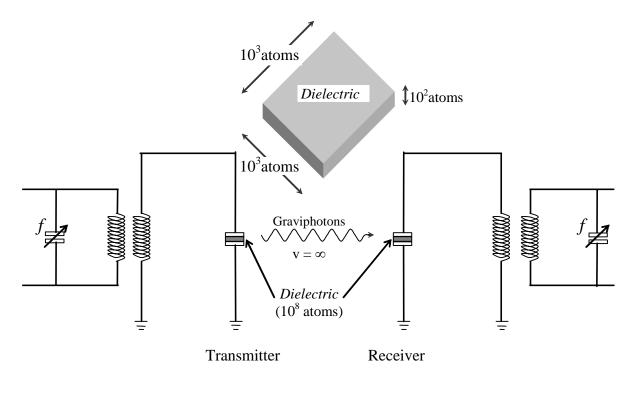


Fig. 16 - Transmitter and Receiver of Virtual Gravitational Radiation.



(a)

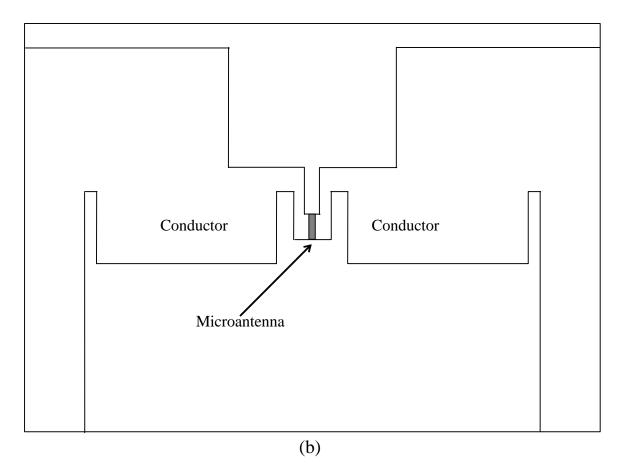


Fig. 17 – Quantum Gravitational Microantenna

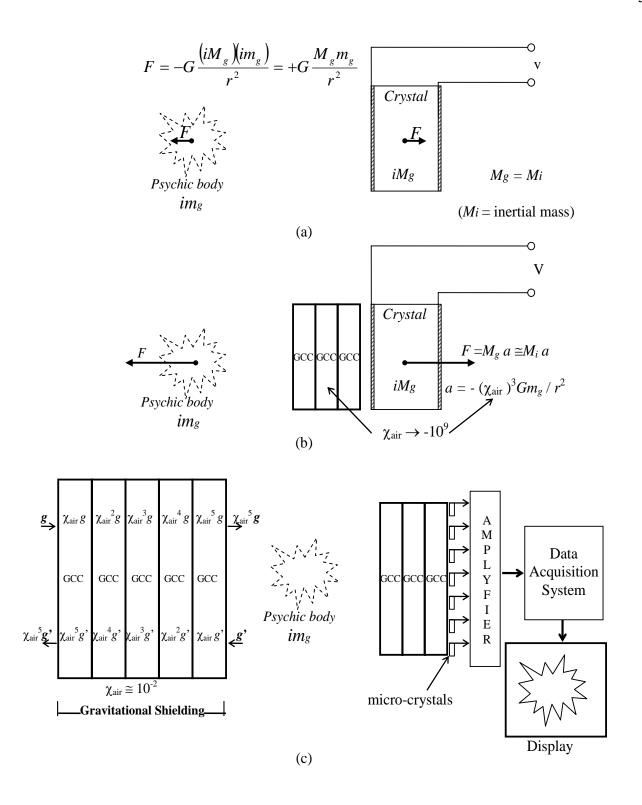


Fig. 18 – Method and device using GCCs for obtaining *images* of *imaginary bodies*.

## REFERENCES

- [1] De Aquino, F. (2006) Mathematical Foundations of the Relativistic Theory Quantum *Gravity*, preprint, physics/0212033.
- [2] Freire, G. F. O and Diniz, A. B. (1973) Ondas Eletromagnéticas, Ed. USP,p.26.
- [3] Halliday, D. and Resnick, R. (1968) Physics, J. Willey & Sons, Portuguese Version, Ed. USP, p.1118.
- [4] Quevedo, C. P. (1977) Eletromagnetismo, McGraw-Hill, p.255 and 269.
- [5] GE Technical Publications (2007), 80044 - F20T12/C50/ECO, GE Ecolux ®T12.
- [6] Aplin, K. L. (2000) PhD thesis, The University of Reading, UK
- [7] Aplin K. L (2005) Rev. Sci. Instrum. 76, 104501.
- [8] Beiser, A. (1967) Concepts of Modern Physics, McGraw-Hill, Portuguese version (1969) Ed. Polígono, S.Paulo, p.362-363.
- [9] Hayt, W. H. (1974), Engineering Electromagnetics, McGraw-Hill. Portuguese version (1978) Ed. Livros Técnicos e Científicos Editora S.A, RJ, Brasil. P.146.
- [10] Benjegerdes, R. et al., (2001) Proceedings of the 2001Particle Accelerator Conference, Chicago. http://epaper.kek.jp/p01/PAPERS/TOAB009.PDF
- [11] Gourlay, S. et al., (2000) Fabrication and Test of a 14T, Superconducting Racetrack Dipole Magnetic, IEE Trans on Applied Superconductivity, p.294.
- [12] BPE soft, Extreme High Altitude Conditions Calculator. http://bpesoft.com/s/wleizero/xhac/?M=p
- [13] Handbook of Chemistry and Physics, 77<sup>th</sup> ed.1996.